



Record-breaking Coronal Magnetic Field in Solar Active Region 12673

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Abstract

The strongest magnetic fields on the Sun are routinely detected at dark sunspots. The magnitude of the field is typically about 3000 G, with only a few exceptions that reported the magnetic field in excess of 5000 G. Given that the magnetic field decreases with height in the solar atmosphere, no coronal magnetic field above ~ 2000 G has ever been reported. Here, we present imaging microwave observations of anomalously strong magnetic field of about 4000 G at the base of the corona in solar active region NOAA 12673 on 2017 September 6. Combining the photospheric vector measurements of the magnetic field and the coronal probing, we created and validated a nonlinear force-free field coronal model, with which we quantify the record-breaking coronal magnetic field at various coronal heights.

Key words: Sun: corona – Sun: magnetic fields – Sun: radio radiation

Supporting material: animation

1. Introduction

The strongest large-scale magnetic field on the Sun is concentrated at certain areas called active regions (ARs). The magnitude of the field is typically about 3000 G, with only a few exceptions that reported the magnetic field in excess of 5000 G (Zirin & Wang 1993; Livingston et al. 2006; Jaeggli 2016; Okamoto & Sakurai 2018; Wang et al. 2018). In their rigorous survey of historical records covering almost 90 years of measuring photospheric magnetic field in ARs, Livingston et al. (2006) reported 55 ARs with magnetic field above 4000 G, five cases above 5000 G, and even one case above 6000 G. In particular, AR 7378 had a magnetic field of $\sim 6,100$ G on 1942 February 28. This AR was associated with a strong geomagnetic storm and gave rise to the discovery of sporadic solar radio emission in the meter-wave band (Livingston et al. 2006). A strong magnetic field is needed to drive extreme solar events; however, given the rarity of the strongest field cases, no reliable association rate has been established.

ARs are called such because they control most of the solar eruptive activity such as solar flares or coronal mass ejections. This activity is powered by the coronal magnetic field, whose energy dominates other forms of energy in the corona. Thus, it is fundamentally important to carefully quantify the coronal magnetic field in ARs. However, no routine diagnostics of the coronal magnetic field at different heights are currently available, so various modeling approaches are used instead.

An exception is the gyroresonant (GR) probing of the coronal magnetic field, which is performed with microwave imaging instruments. The foundation of this method is very simple. The coronal plasma is optically thick at a few of the lowest harmonics of the local gyro-frequency. The brightness temperature of this optically thick emission is equal to the kinetic temperature of the emitting plasma. The stronger the local magnetic field, the higher the radio frequency at which the given source is bright. For a given line of sight, the magnetic field increases toward the solar surface, while the temperature decreases sharply at the transition region between the corona and

the chromosphere. Thus, the radio brightness will drop sharply at the emission frequency that corresponds to the gyroresonant frequency at the transition region. Investigating the radio brightness pixel by pixel, we can recover the magnetic field strength at the base of the corona, provided that broadband microwave data with high spatial resolution are available.

In practice, these diagnostics are often available either at a few single frequencies or over a limited spectral range. For a typical AR, the GR emission dominates below ~ 10 GHz, while stronger ARs show bright GR emission at 17–18 GHz (Lee 2007; Shibasaki et al. 2011). No GR emission at higher frequencies has ever been reported, implying that the coronal magnetic field does not typically exceed ~ 2100 G.

However, AR 12673 demonstrated an extremely strong photospheric field of about 5000 G over a rather extended area of the sunspot during the day of 2017 September 6. In addition to this, a light bridge at the polarity inversion line has occasionally demonstrated extreme values of the transverse magnetic field up to 5570 G (Wang et al. 2018). Magnetic field of such magnitude was formed in association with an unusually fast new magnetic flux emergence (Sun & Norton 2017). During the period of 2017 September 3–4, a number of magnetic bipoles with a scattered structure subsequently emerged on the eastern side of the pre-existing sunspot. Their complex mutual motion leads to the formation of the high sheared strong magnetic configuration that eventually produced an X9.3 flare on 2017 September 6. A comprehensive description of the photospheric magnetic field evolution of AR 12673 can be found in Yang et al. (2017). Although this strongest transverse magnetic field may close very low at the chromosphere and have no appreciable fingerprint in the corona, the more extended magnetic field of the order of 5000 G may reach the higher coronal heights. If this is the case, a bright GR coronal source might appear at an unusually high microwave frequency.

Here, we report on such a bright GR source at the highest available frequency, 34 GHz, using Nobeyama Radio Heliograph (NoRH) data during the day of 2017 September 6, which is twice larger than the largest frequency at which a

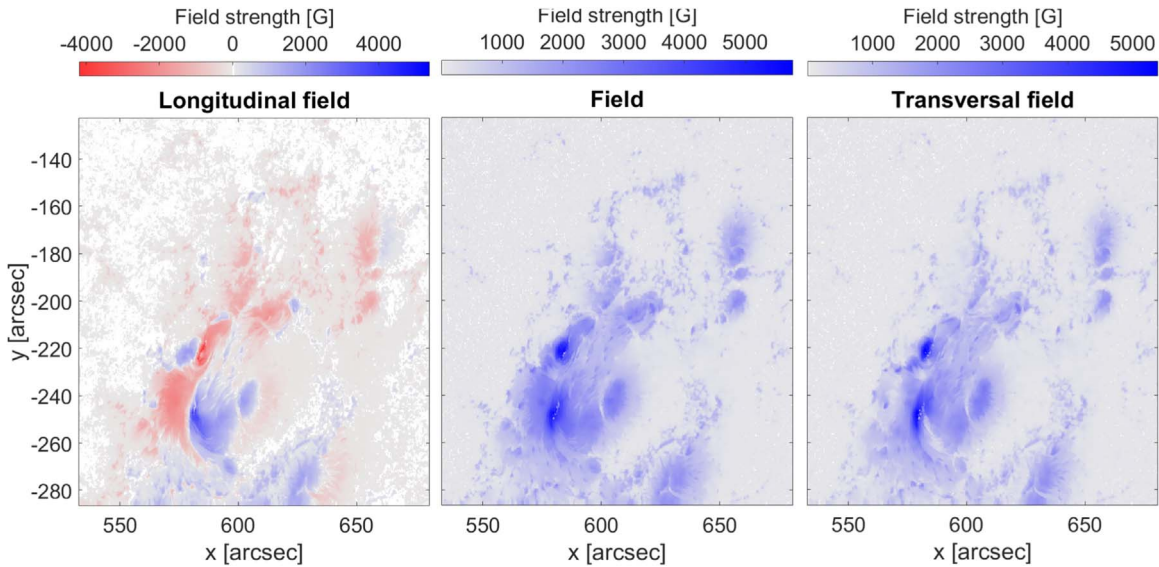


Figure 1. Map of the absolute value of the magnetic field (central panel), its longitudinal (left panel) and transversal (right panel) components observed by *Hinode* SOT/SP at 18:36:50 UT on 2017 September 6.

GR source has ever been reported. This observation alone unambiguously indicates the presence of the coronal magnetic field of up to 4000 G, which is twice as large as values reported thus far.

2. Observations and Data Analysis

2.1. Overview

In this study we employ spectropolarimetry data obtained with the Helioseismic and Magnetic Imager Instrument (HMI) on board the *Solar Dynamics Observatory* (SDO; Pesnell et al. 2012; Schou et al. 2012) and *Hinode* Solar Optical Telescope/Spectro-Polarimeter (SOT/SP; Kosugi et al. 2007; Tsuneta et al. 2008) to derive the photospheric vector magnetograms needed to initiate coronal magnetic models as well as microwave images obtained with the Nobeyama Radio Heliograph (NoRH; Nakajima et al. 1994) to quantify the coronal magnetic field.

Given that AR 12673 demonstrated remarkably unusual behavior, many standard pipeline products contain artifacts, while standard analysis techniques required manual corrections as described below. For our analysis, we focused on two time frames: 03:34:42 UT and at 18:36:50 UT on 2017 September 6. The first time frame is selected during local mid-day of the microwave observations at the Nobeyama Observatory that detected a unique long-living GR source at 34 GHz, while the amount of SDO/HMI artifacts was minimal during that time period. The second time frame is selected based on the time when the largest photospheric magnetic field was detected by *Hinode* SOT/SP.

2.2. Inversion of the *Hinode* SOT/SP Stokes Profiles

NOAA AR 12673 has been observed by both SDO/HMI and *Hinode* SOT/SP instruments. In the standard vector magnetograms available at <http://jsoc.stanford.edu/> for SDO/HMI and at https://csac.hao.ucar.edu/sp_data.php for *Hinode* SOT/SP, the magnetic field is artificially bounded below 5000 G. For this AR, the magnetic field is saturated at this threshold.

To measure the real (unsaturated) value of the magnetic field from Stokes profiles observed by *Hinode* SOT/SP at 18:36:50

UT on 2017 September 6, we use the VFISV inversion code (Borrero et al. 2011) available at <https://www2.hao.ucar.edu/csac/csac-spectral-line-inversions>. We modified this source code to allow magnetic fields above 5000 G and applied to the measured Stokes profiles. This way we obtained an unsaturated vector magnetogram. The corresponding maps of the absolute value, as well as longitudinal and transversal components of the magnetic field vector, are shown in Figure 1.

2.3. SDO/HMI Vector Magnetograms: Removing Inversion Artifacts

Standard SDO/HMI vector magnetograms obtained on 2017 September 6 contain apparent artifacts in several pixels, which are most likely produced by an extremely fast emergence of the magnetic flux (Sun & Norton 2017). These artifacts (glitches) are obvious from the bottom-left panel of Figure 2 (suspicious pixels are marked with green frames): both LOS and transversal components are totally different from most of their neighborhoods. This is more evident for the transversal component. We manually eliminated these glitches as follows: we identified a “bad” pixel with the largest number of “good” neighboring pixels and replaced the magnetic field components of this pixel with the average of its “good” neighbors. Then we applied this algorithm recursively to all suspicious pixels. The result of such procedure is presented in the bottom-right panel of Figure 2.

2.4. Performing π -disambiguation

Vector magnetic field inverted from the Stokes profiles is intrinsically ambiguous. The direction of the transversal component can be measured only up to 180° accuracy, meaning that for every pixel in the magnetogram there are two possible directions of the transversal magnetic field. Thus, one has to perform a π -disambiguation procedure before using the magnetogram as a photospheric boundary condition for nonlinear force-free field (NLFFF) reconstruction. In this work, we employ two disambiguation methods: SFQ (Rudenko & Anfinogentov 2014) and the Minimum Energy method (Metcalf 1994). Both methods give almost identical results in

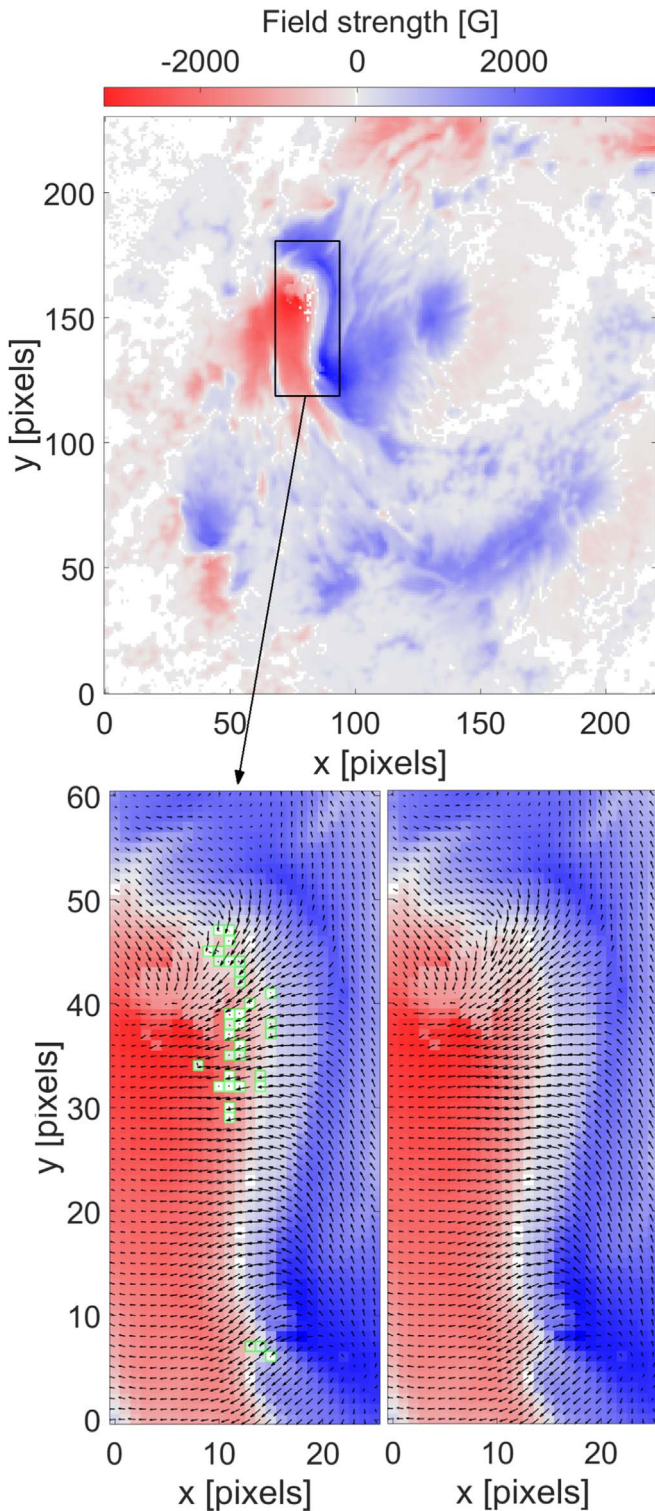


Figure 2. Magnetic field of the AR 12673 provided by *SDO*/HMI at 2017 September 6 03:34:42. Top row: Overview of the magnetic field (the pixel's color represents the value of the LOS-component). Bottom row, left: selected field of view (FOV) with glitches (marked with green frames). Bottom row, right: the same FOV with removed glitches. Black arrows show direction and relative strength of the transversal component of the magnetic field.

the strong field areas, which validates the π -disambiguations performed.

In the case of the magnetograms taken at 18:36:50 UT, there is an evident disambiguation artifact present in the outputs of both disambiguation methods as well as in the standard *SDO*/HMI

data product. To remove this artifact, we manually fix the field direction in the corrupted area and run the π -disambiguation procedures under this constraint. This successfully removes the artifact.

2.5. Maps of the GR Source at 34 GHz

Standard data products of NoRH include daily images at 17 and 34 GHz and 10 minutes images for 17 GHz. To produce microwave images at 34 GHz with a higher cadence (10 minutes) than available from the NoRH pipeline, we use the software developed by the NoRH team and available at <https://solar.nro.nao.ac.jp/norh/archive.html>. The image reconstruction has been performed using Hanaoka synthesis program that implements the classical CLEAN algorithm (Högbom 1974).

These images consistently show high brightness temperature indicative of the GR emission process at least up to 34 GHz. The centroid location of the 34 GHz images, however, noticeably fluctuates from image to image within $\sim 5''$ in a random direction. In Figure 3 we shifted the 34 GHz image by $3''$ to the south, which is within the determined NoRH positioning accuracy.

2.6. NLFFF Reconstruction of Coronal Magnetic field

The photospheric vector magnetograms, with fixed artifacts and carefully resolved azimuth ambiguity, were prepared for two time frames of interest, specifically, 03:34:42 UT and 18:36:50 UT on 2017 September 6. These magnetograms were employed as the bottom boundary conditions to perform NLFFF reconstruction of the coronal magnetic field. Specifically, we employed two different versions of the optimization method (Wheatland et al. 2000; Wiegelmann 2004) developed and validated by our team (Fleishman et al. 2017) using a full-fledged 3D magnetohydrodynamic model as a proxy of a real AR. Fleishman et al. (2017) found that existing methods of the photospheric magnetic boundary condition preprocessing, proposed to remove the forced component of the magnetic field, in fact result in a corrupted height scale, and do not improve the quality of the reconstruction. Therefore, we did not apply any preprocessing, but used the bottom boundary condition as is. Having unusually strong photospheric magnetic field implies a proportionally lower plasma beta, which further justifies the use of the force-free approximation. In what follows, we will validate our NLFFF reconstructions by comparison of the modeled values with those derived from the radio observations at the base of corona, where the force-free approximation holds much better. Both reconstructions resulted in very similar coronal magnetic structures, which provides confidence in the obtained solutions. To avoid any confusion, we use here only one of the solutions; namely, the one obtained with the weighted optimization code (Fleishman et al. 2017).

2.7. Chromospheric and Coronal Thermal Model

The presence of a strong magnetic field of ~ 4 kG at the coronal heights, where the thermal plasma is hot, is an absolute prerequisite of having a bright GR radio source at 34 GHz, like the one reported in Section 2.5. The magnetic model alone, however, is insufficient to compute the radio brightness needed for a meaningful model-to-data comparison, which in addition requires a thermal model on top of the magnetic one. Thus, to compute simulated radio emission from the model, its volume

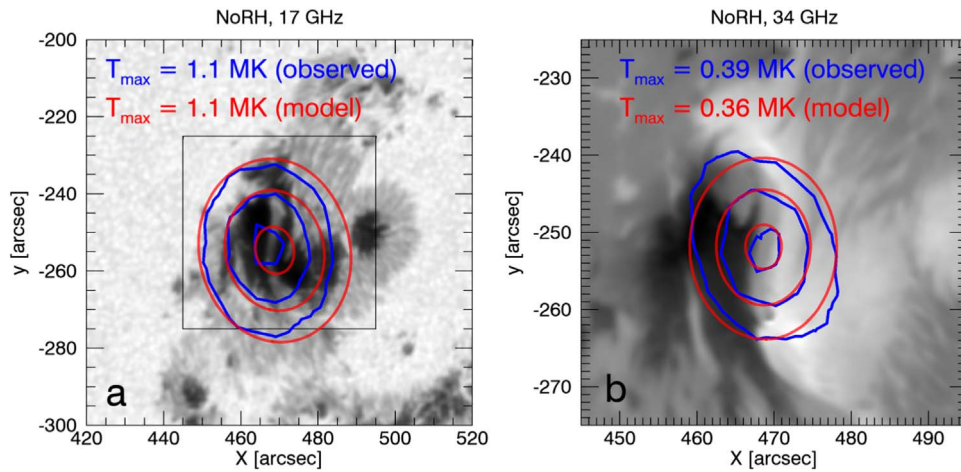


Figure 3. Continuum white light emission (a) and LOS magnetogram (b) obtained by *SDO*/HMI on 2017 September 6 03:34:42. Blue contours show microwave radio emission observed with NoRH on 03:30:01 UT at 17 GHz (a) and 34 GHz (b). The synthetic radio maps calculated from our 3D magneto-thermal model are shown by the red contours. In all cases, the contours are plotted at the levels of 20%, 50%, and 90% of the peak brightness temperature. The peak brightness temperatures are given by the labels in the top-left corner of each panel. The black rectangle shown on the left panel indicates the FOV of the panel (b). Note that the radio sources project on the light bridge (a bright vertical feature seen in panel (a) at the center of the radio source), which is elongated along the polarity inversion line seen in the magnetic field LOS map shown in panel (b).

has to be filled in with a radiating thermal plasma. Within the GX Simulator methodology (Nita et al. 2018), this process is called the *volume reprocessing*. During the reprocessing, the length and the mean value of the magnetic field along the field line associated with a given volume element (voxel) are computed for each voxel of the 3D magnetic model. Then, a parametric model of the coronal plasma heating is applied to every voxel crossed by a closed line of the magnetic field obtained from the reconstruction. A thermal model of the chromosphere, which has the nonuniform adaptive height spacing needed to resolve the transition region, is added at the bottom part of the datacube following a photospheric mask obtained from a joint analysis of the white light and LOS magnetic field maps (Nita et al. 2018). This way, the magnetic “skeleton” is getting filled with a plasma with appropriate temperature and density, forming a thermal structure of our model.

3. Discussion

Starting from the photospheric vector magnetograms, we built 3D coronal magnetic models using NLFFF reconstruction, populated these datacubes with coronal and chromospheric thermal plasma (Nita et al. 2018), and validated them by comparison of the simulated microwave emission with the microwave imaging data at 17 and 34 GHz. An almost perfect match between the simulated and observed images (see Figure 3) validates both magnetic and thermal structure of our 3D model.

Figure 4(a) shows a 3D representation of the coronal magnetic field by displaying isogauss surfaces of 3500 G (green wired surface), 2000 G (red wired surface), and 1000 G (blue wired surface). This figure demonstrates that the kilogauss magnetic field occupies a highly significant coronal volume at the AR, extending to much higher heights than in a typical case. Figure 4(b) gives a complementary view of the magnetic structure by showing a flux tube of the strongest, highly twisted magnetic field forming a flux rope in the core of the AR.

To quantify the height dependence of the coronal magnetic field, Figure 5 displays the dependence of the maximum magnetic field value at a given height (the green line) along with the estimate for the magnetic field at the base of the corona derived from the GR emission at 34 GHz. We also plot the historically largest measurements of the magnetic field obtained by Brosius & White (2006) at two different heights from Very Large Array (VLA) observations above the limb. It is interesting that the upper bounds of these measurements match our model surprisingly well. This indicates that AR 12673 did show the strongest, record-breaking coronal magnetic field. However, given the proximity of the VLA data points to our curve that displays the largest magnetic field versus height, the comparably strong coronal magnetic fields might be more common than has been appreciated so far. To check this expectation, we looked at the historical daily record of the NoRH data at 34 GHz and indeed found many cases showing unexpectedly bright radio sources indicative of extremely strong coronal field. Some of these cases correspond to the strongest sunspot magnetic field reported by Okamoto & Sakurai (2018). A detailed analysis of this data set is underway and will be published elsewhere shortly.

Wang et al. (2018) reported strong photospheric magnetic field of 5570 G at AR 12673 based on Goode Solar Telescope data, detected roughly half a day after our observations of the bright GR source. To investigate whether this enhancement of the photospheric magnetic field had any effect on the coronal magnetic field, we obtained the photospheric vector magnetic field using *Hinode* SOT/SP data from a restricted FOV at 18:36:50 UT that demonstrated a consistently strong magnetic field up to 5700 G at the light bridge, embedded this magnetogram into a larger FOV *SDO*/HMI magnetogram, and produced an NLFFF reconstruction of the coronal magnetic field. Then, we computed the largest magnetic field at each height and plotted it in Figure 5 in red. Interestingly, the red curve exceeds the level of the green one (obtained for 03:34:42 UT) only over a very restricted range of the heights low in the solar atmosphere, while then merges the green one.

This finding suggests that the strong transverse photospheric magnetic field, reported by Wang et al. (2018), closes at the

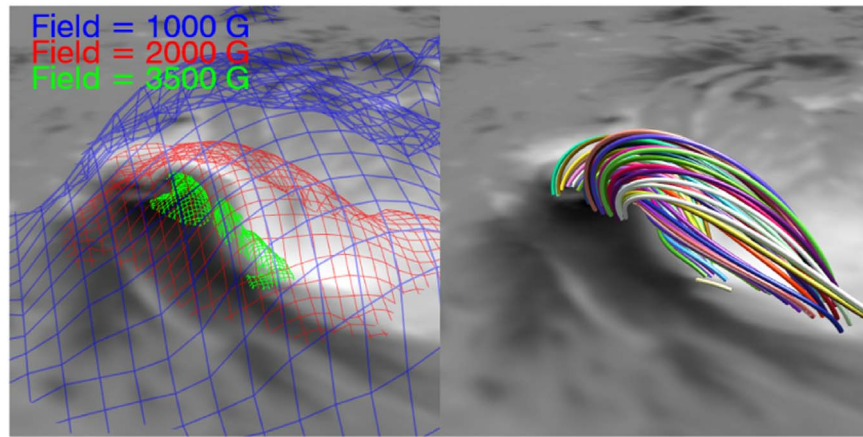


Figure 4. Isogauss layers corresponding to the magnetic field strength of 1000 G (blue), 2000 G (red), and 3500 G (green). The right panel shows the magnetic field lines rooting at the photospheric areas where the magnetic field exceeds 3500 G.

(An animation of this figure is available.)

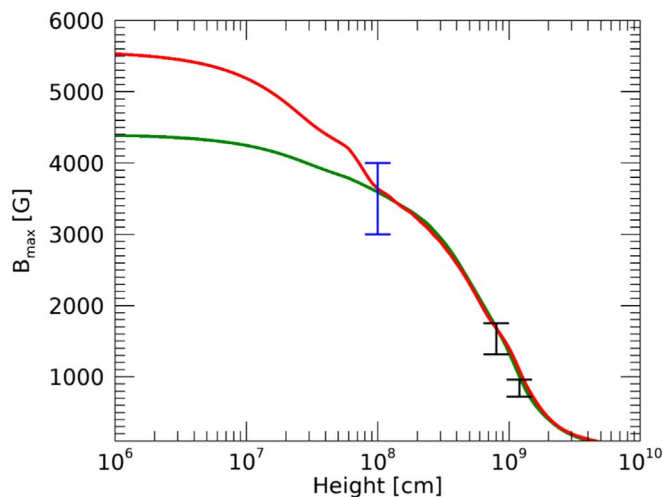


Figure 5. Height dependence of the strongest magnetic field obtained from the NLFFF models for 03:36 UT (green line) and 18:36 UT (red line). Magnetic field estimate at the base of the corona obtained from NoRH data is shown with a blue vertical dash. The strongest coronal magnetic field ever reported (Brosius & White 2006) based on VLA observations above the limb is shown with black vertical dashes.

photosphere or low chromosphere, but does not propagate upward to the corona. In contrast, a bigger area of a slightly weaker, but still very strong, magnetic field of about 5000 G, which persists in this AR during the day of September 6 or even longer, has a significant imprint in the coronal magnetic field.

4. Conclusions

It is extremely exciting that AR 12673 shows a unique, long-living bright GR source at 34 GHz (see Figure 3, right panel), which is twice as large than the largest frequency at which a GR source has ever been reported and indicative of an unexpectedly strong coronal magnetic field. In addition, we found that the historically strongest coronal magnetic field values, obtained by Brosius & White (2006) at two different heights from VLA observations above the limb, match the obtained here height dependence of the strongest magnetic field remarkably well. This implies that strong coronal magnetic

fields, comparable to that reported here, might be a far more common phenomenon than has been appreciated so far.

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Erratum: “Record-breaking Coronal Magnetic Field in Solar Active Region 12673” (2019, ApJL, 880, L29)

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Figure 2 in the published article contains an error: the transverse components B_x and B_y of the magnetic field shown with arrows in the lower panel of Figure 2 were misinterpreted, respectively, as B_y and B_x by mistake. Only the figure is affected by this error. The corrected Figure 2 is presented here.

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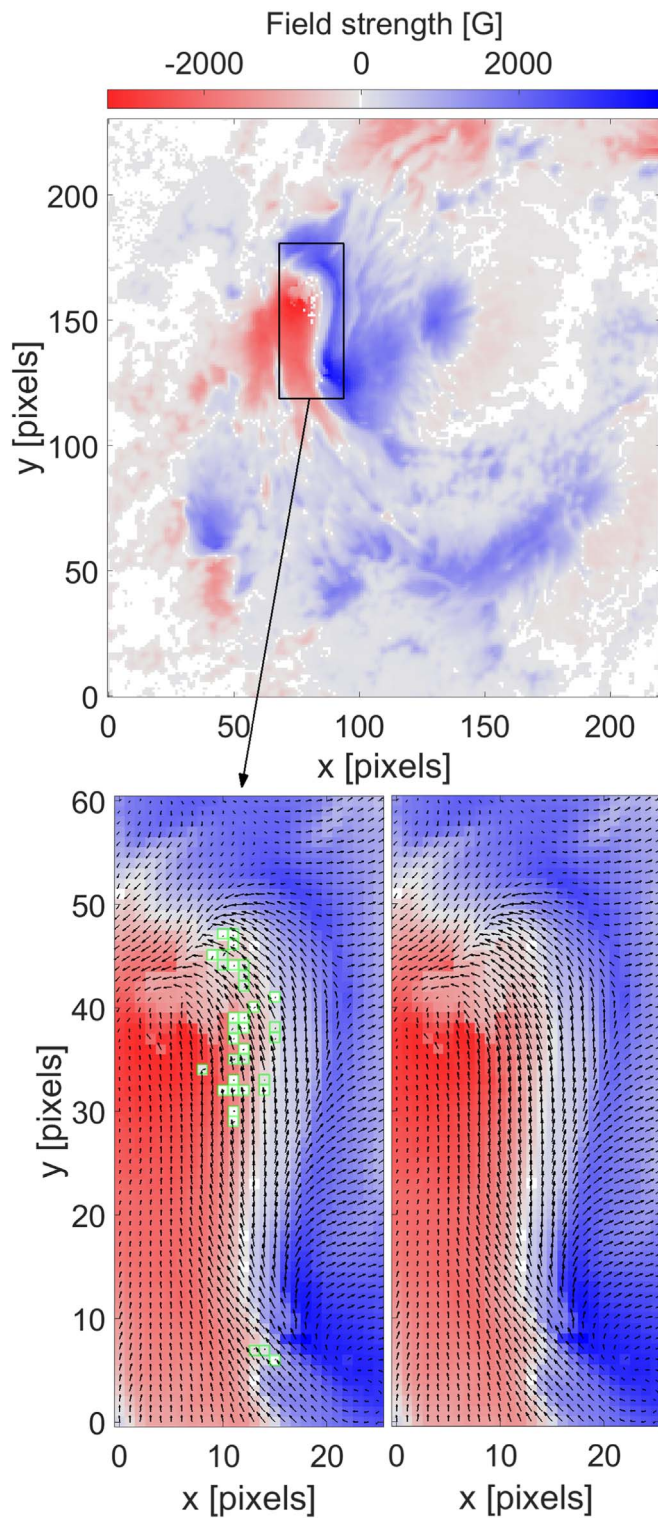


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